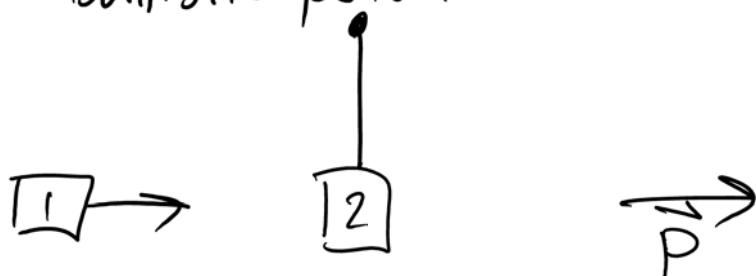


pdf of Tuesday lecture: ter.ps/phys410sep04

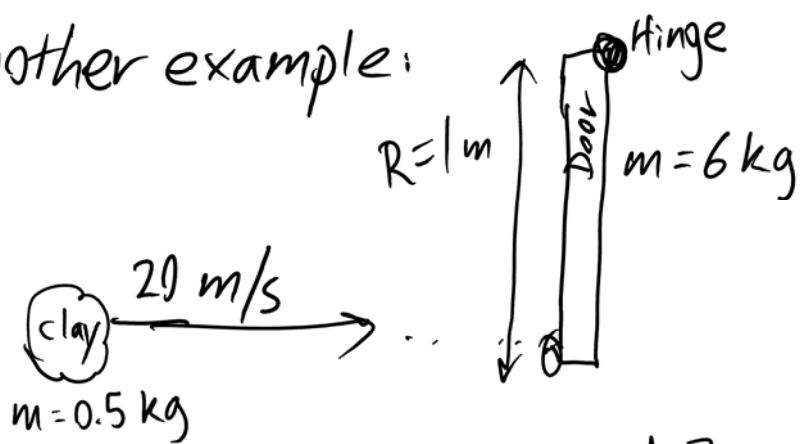
From today's lecture: ter.ps/phys410sep06

What's conserved in a collision?

Last time: Ballistic pendulum



Another example:



Clay sticks to door

Q: How fast will the door swing after it's hit?

what will be conserved? Not energy.

fixed pivot \Rightarrow angular momentum will be conserved

Angular momentum (around the hinge)

$$\text{Before collision: } \vec{L} = \vec{r} \times \vec{p} \Rightarrow |\vec{L}| = r p \sin \theta$$

$$= r_{\perp} p$$

$$= (1 \text{ m}) \times (10 \text{ kg m/s}) \\ = 10 \text{ kg m}^2/\text{s}$$

After collision: $L = I_{\text{total}} \omega$

$$I_{\text{total}} = I_{\text{clay}} + I_{\text{door}}$$

$$= m_{\text{clay}} R^2 + \frac{1}{3} m_{\text{door}} R^2 = 2.5 \text{ kg}\cdot\text{m}^2$$

Aside: $I_{\text{door}} = \int dm r^2$ 

$$= \int_0^R \left(\frac{dr}{R} m_{\text{door}} \right) r^2$$

$$= \frac{m_{\text{door}}}{R} \frac{1}{3} R^3$$

$$\Rightarrow \omega = \frac{L}{I_{\text{total}}} = \frac{10 \text{ kg}\cdot\text{m}^2/\text{s}}{2.5 \text{ kg}\cdot\text{m}^2} = 4 \text{ s}^{-1} = 4 \text{ rad/s}$$

Check: Was linear momentum conserved?

Before: $P = 10 \text{ kg}\cdot\text{m/s}$

After: $P = P_{\text{clay}} + P_{\text{door CM}}$

$$= m_{\text{clay}}(WR) + m_{\text{door}}\left(\omega \frac{R}{2}\right)$$

$$= (0.5 \text{ kg})(4 \text{ rad/s})(1 \text{ m}) + (6 \text{ kg})(4 \text{ rad/s})(0.5 \text{ m})$$

$$= 2 + 12 = 14 \text{ kg}\cdot\text{m/s}$$

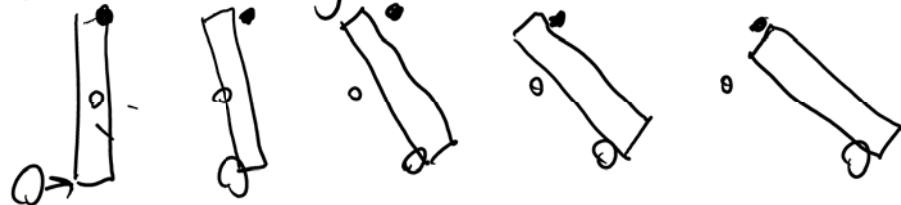
$\Rightarrow \vec{P}$ is not conserved in this collision

$$\Delta \vec{P} = \vec{P}_f - \vec{P}_i = 4 \text{ kg}\cdot\text{m/s} \quad \times$$

So the impulse was $4 \text{ kg}\cdot\text{m/s}$ \times

$$= \int \vec{F} dt$$

If there had been no hinge:



$$P_f = 10 \text{ kg}\cdot\text{m/s} \text{ in this case}$$

Bat demo:

Center of percussion

Hamilton's principle

The path that a dynamical system actually takes from one point to another over a given time interval is one that makes the action integral stationary.

Action integral : $S = \int_{t_1}^{t_2} \mathcal{L} dt$

Path: $x(t)$, i.e. position vs. time

points: starting and ending points of the path
 $(t_1, x_1) \rightsquigarrow (t_2, x_2)$
with the endpoints fixed

Note: S is a scalar, calculated from a function
(the path) that can vary

So this action integral is a "functional"

Stationary: small variations in the path do not change S to first order

{minimum, maximum, inflection point} \rightarrow stationary

\rightarrow Solving this with the right \mathcal{L} will tell us how the physical system behaves!

Use this... find the path which makes it stationary

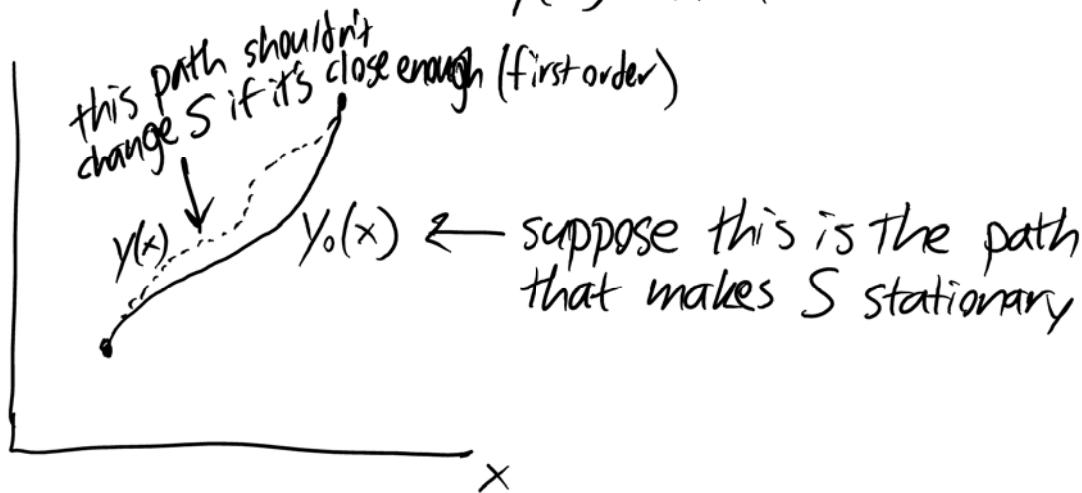
Calc. of variations

We have $S = \int_{x_1}^{x_2} f(y, y', x) dx$

$y(x)$ is the path

$$y' = \frac{dy}{dx}$$

we want to find the $y(x)$ which makes this stationary



write as $y(x) = y_0(x) + \alpha \eta(x)$

$\eta(x)$ is not small
but $\alpha \rightarrow 0$

How does S change from this?

$$\begin{aligned} \frac{\delta S}{\delta \alpha} &= \int_{x_1}^{x_2} \frac{\partial f}{\partial \alpha} dx = \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial y} \frac{\partial y}{\partial \alpha} + \frac{\partial f}{\partial y'} \frac{\partial y'}{\partial \alpha} + \frac{\partial f}{\partial x} \frac{\partial x}{\partial \alpha} \right) dx \\ &= \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial y} \eta(x) + \frac{\partial f}{\partial y'} \eta'(x) \right) dx \end{aligned}$$

use I.B.P. $\int v du = uv - \int u dv$

$$\int_{x_1}^{x_2} \underbrace{\frac{\partial f}{\partial y'}}_{\text{"V"}}, \underbrace{\eta'(x) dx}_{\text{"dU"}}, = \left[\eta(x) \frac{\partial f}{\partial y'} \right] \Big|_{x_1}^{x_2} - \underbrace{\int_{x_1}^{x_2} \eta(x) \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) dx}_{=0 \text{ since } \eta(x_1) = \eta(x_2) = 0}$$

$$\Rightarrow \frac{ds}{dx} = \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \right] \eta(x) dx$$

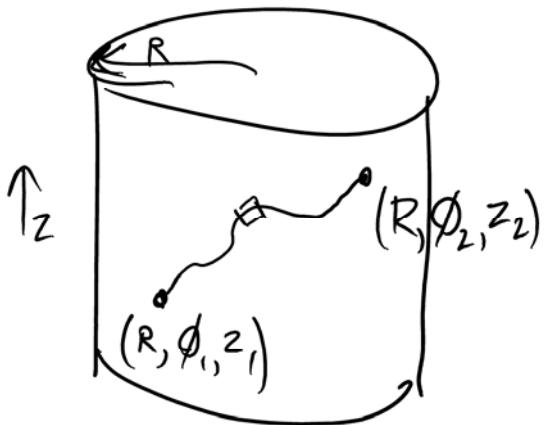
want this to be zero for any $\eta(x)$

$$\Rightarrow \frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0$$

This is the Euler-Lagrange equation

Means: The $y(x)$ [path] which satisfies this equation makes S stationary

Example: Shortest path between two points on a cylinder



Length of small piece

$$ds = \sqrt{dz^2 + (R d\phi)^2}$$

suppose that ϕ is a function of z

$$\Rightarrow ds = \sqrt{1 + R^2 (\frac{d\phi}{dz})^2} dz$$

$$l(R, \phi_1, \zeta_1) \rightarrow ds = \sqrt{1 + R^2 \left(\frac{d\phi}{dz}\right)^2} dz$$

$$\text{Total length: } L = \int ds = \int_{z_1}^{z_2} \sqrt{1 + R^2 \phi'^2} dz$$

want to minimize L by varying $\phi(z)$

$$\text{NSR E-L equation with } f = \sqrt{1 + R^2 \phi'^2}$$

$$\frac{\partial F}{\partial \phi} = \frac{d}{dz} \left(\frac{\partial F}{\partial \phi'} \right)$$

$$0 = \frac{d}{dz} \left(\frac{1}{2} (1 + R^2 \phi'^2)^{1/2} \cdot R^2 \cdot 2\phi' \right)$$

$$\Rightarrow \frac{R^2 \phi'}{\sqrt{1 + R^2 \phi'^2}} = C$$

$$R^4 \phi'^2 = C^2 (1 + R^2 \phi'^2) \rightarrow \phi'^2 = \frac{C^2}{R^4 - C^2 R^2}$$

a constant
call that D

$$\Rightarrow \phi'(z) = \sqrt{D}, \text{ a constant}$$

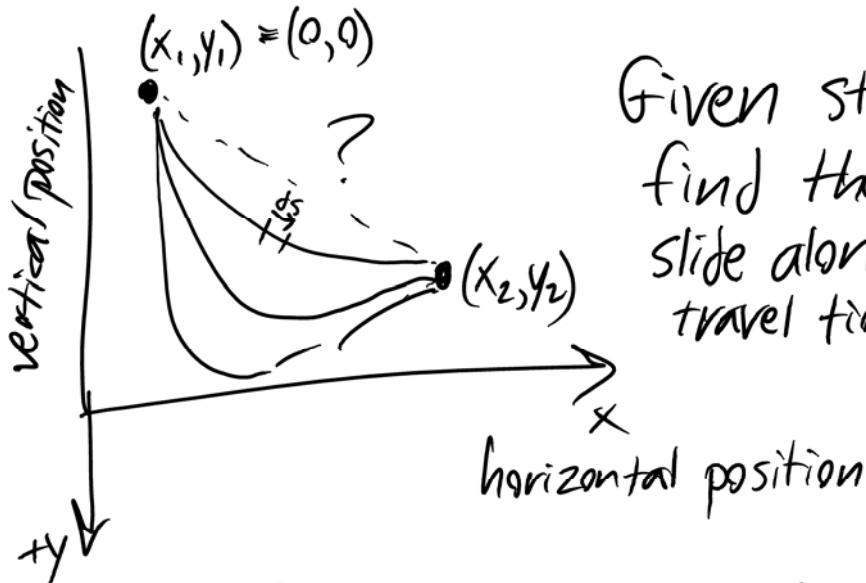
$$\Rightarrow \phi(z) = \phi_0 + \sqrt{D} z$$

$\uparrow \quad \uparrow$
Impose fixed endpoints

This was for a 1-D path, i.e. $y(x)$ or $\phi(z)$

In 2 or more dimensions, get an E-L equation
for each of the coordinates

Brachistochrone



Given starting & ending points,
find the track for a object to
slide along that minimizes the
travel time

Describe that path $x(y)$

"S" functional is the total time

$$t = \int_{\text{path}} dt = \int_{\text{path}} \frac{ds}{v} \quad \text{speed there}$$

use Energy conservation:

$$mgy = \frac{1}{2}mv^2 \Rightarrow v = \sqrt{2gy}$$

$$ds = \sqrt{dx^2 + dy^2} = \sqrt{\left(\frac{dx}{dy}\right)^2 + 1} dy$$

$$t = \int \frac{\sqrt{x'^2 + 1}}{\sqrt{2gy}} dy$$

$\underbrace{}$
this is "f(x, x', y)"

use E-L equation:

$$\frac{\partial f}{\partial x} = \frac{d}{dy} \left(\frac{\partial f}{\partial x'} \right)$$

"Q" $\Rightarrow \frac{\partial f}{\partial x}$, is constant = C

$$\Rightarrow \frac{\frac{1}{2}(x'^2+1)^{-1/2} \cdot 2x'}{\sqrt{2gy}} = C \quad \Rightarrow \frac{x'}{\sqrt{x'^2+1}} = \underbrace{\sqrt{2gC}}_{C'} \sqrt{y}$$

$$\Rightarrow x'^2 = C'^2 y (x'^2 + 1)$$

$$x'^2(1 - C'^2 y) = C'^2 y \quad \Rightarrow \quad x' = \sqrt{\frac{C'^2 y}{1 - C'^2 y}}$$

Need to solve this D.E. to get our $x(y)$

Integrate... how?

$$\text{Taylor: } y = a(1 - \cos \theta) = \frac{1}{2C'^2}(1 - \cos \theta)$$

$$\Rightarrow x' = \sqrt{\frac{\frac{1}{2}(1 - \cos \theta)}{1 - \frac{1}{2}(1 - \cos \theta)}} = \sqrt{\frac{1 - \cos \theta}{1 + \cos \theta}}$$

$$\text{Mult. by } \sqrt{\frac{1 - \cos \theta}{1 - \cos \theta}} \Rightarrow = \frac{1 - \cos \theta}{\sin \theta}$$

$$\text{Also } dy = \frac{1}{2C'^2} \sin \theta d\theta$$

$$\Rightarrow x = \int \frac{1 - \cos \theta}{\sin \theta} \cdot \frac{1}{2C'^2} \sin \theta d\theta$$

$$\Rightarrow \text{Get solution } x = \frac{1}{2C'^2} (\theta - \sin \theta) = a(\theta - \sin \theta)$$

$$y = a(1 - \cos\theta)$$

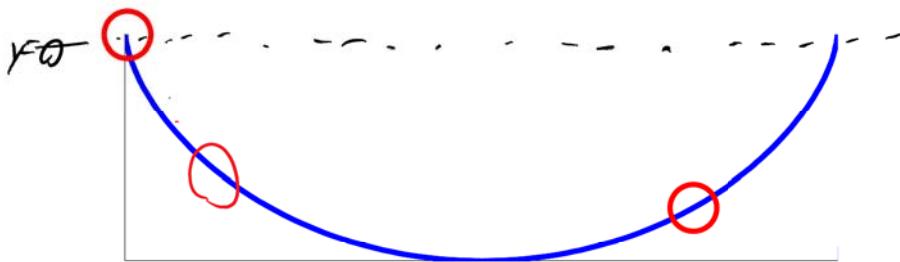
parametric solution with parameter θ
(x no longer has to be a single-valued function of y)

One adjustable parameter, a , which we
must set to satisfy the boundary condition
i.e., ends at (x_2, y_2)

Brachistochrone solution

Thursday, September 06, 2012

7:48 AM



cycloid

There is sometimes a unique solution,
and sometimes multiple solutions, depending on (x_2, y_2)

Comments:

For small θ :

$$y = a(1 - \cos \theta) \sim \frac{a}{2} \theta^2$$

$$x = a(\theta - \sin \theta) \sim \frac{a}{6} \theta^3$$

- Slope $\sim \frac{y}{x}$ is $\sim \frac{1}{3\theta}$

infinite for small θ ,
i.e. it departs vertically
downward at first

- $t = \frac{1}{\sqrt{2g}} \int_0^{y_2} \frac{\sqrt{x'^2 + 1}}{\sqrt{y}} dy$

$$x' = \frac{dx}{dy} = \frac{\frac{dx}{d\theta}}{\frac{dy}{d\theta}} = \frac{1 - \cos \theta}{\sin \theta}$$

$$= \frac{1}{\sqrt{2a}} \left(\int_{\theta_0}^{\theta_F} \sqrt{2a} d\theta \right)$$

$$= \sqrt{\frac{1}{2g}} \int_0^{\pi} \sqrt{2a} d\theta$$

$$= \sqrt{\frac{a}{g}} \theta_f$$